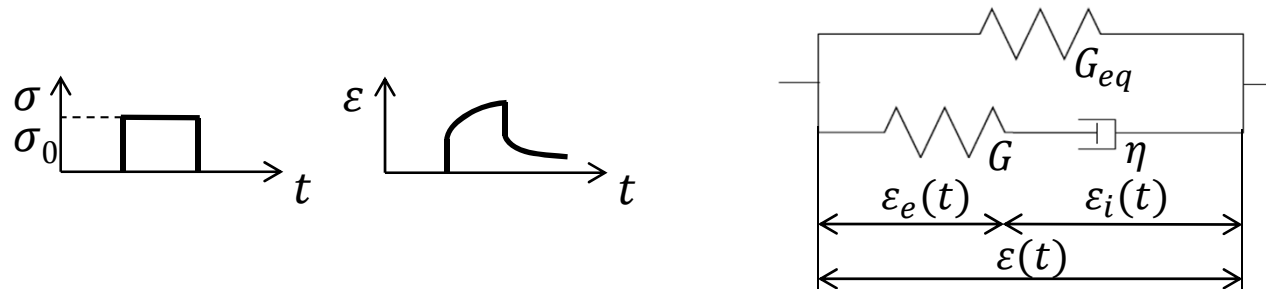


# Modelling of viscoelastic materials with LS-Dyna



11<sup>th</sup> German LS-Dyna Forum 2012, Ulm

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DYNAmore GmbH

Paul Du Bois  
Consultant



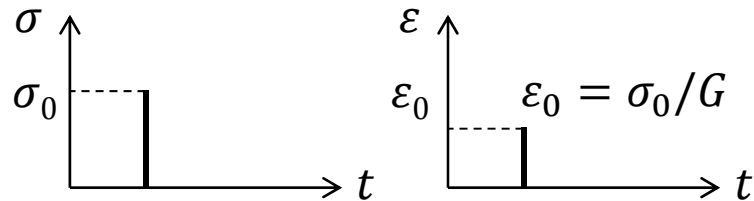
# Overview

- **Motivation**
- **Models for viscoelastic materials in LS-Dyna**
  - Generalized Maxwell Model
  - Tabulated hyperelasticity
- **Modelling of a rubber material**
  - MAT\_SIMPLIFIED\_RUBBER
  - MAT\_OGDEN\_RUBBER
  - BioRID Jacket Certification Test

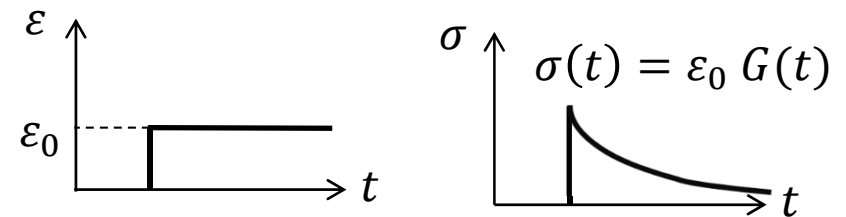
# Motivation

- Characteristic properties of viscoelastic solids

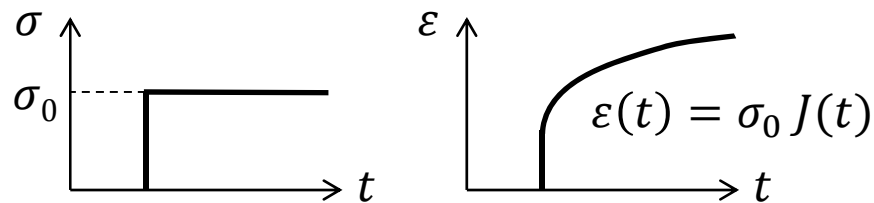
- Instantaneous elasticity



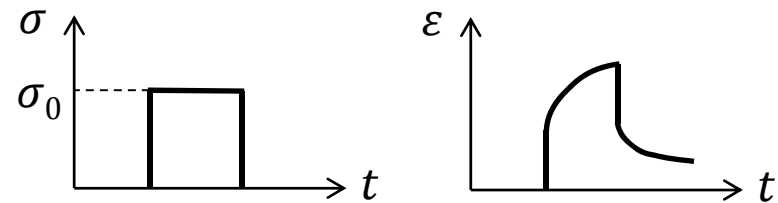
- Stress relaxation under constant strain



- Creep under constant stress



- Instantaneous and delayed recovery



- Many materials show viscoelastic characteristics

- Rubbers
- Foams
- Thermoplastics
- Composites
- ...

# Models for viscoelastic materials in LS-Dyna

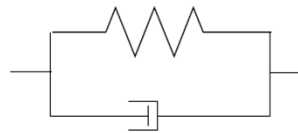
- Linear viscoelastic material models based on rheological models

- Material models: 6, 61, 76, 86, 134, 164, 234, 276,...

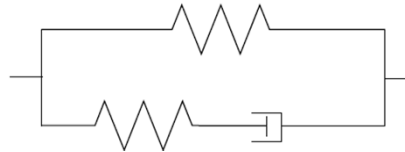
Maxwell-Element



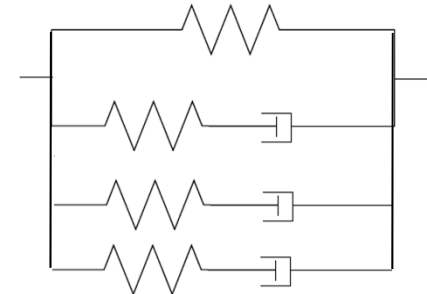
Kelvin-Voigt-Element



Standard linear Solid



Generalized Maxwell Element



- Material model (equilibrium stress) + viscoelastic overstress

- Material models: Hyperelasticity, (Visco-) Plasticity, ...

57, 73, 77, 87, 91, 124, 127, 129, 155, 158, 175, 178,...

- Viscoelastic Overstress: Generalized Maxwell Element

- Material models with elastic and strain rate dependent characteristics

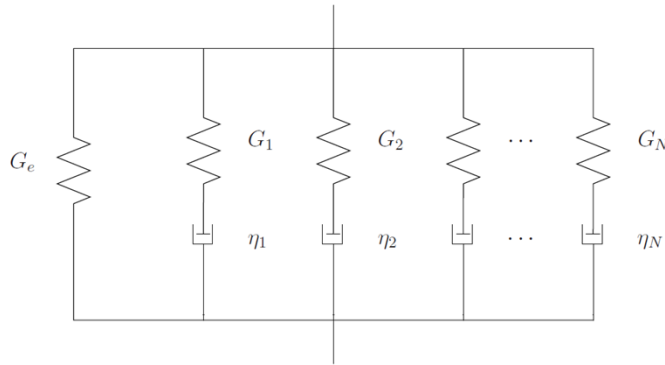
- Rubbers and Foams: MAT\_181 (MAT\_SIMPLIFIED\_RUBBER/FOAM), MAT\_083

- Creep: MAT\_115, MAT\_188

- Quasilinear viscoelasticity: MAT\_176

# Generalized Maxwell Element

- Rheological model with linear elements



Spring (Hooke)  $\sigma(t) = G\varepsilon(t)$

Dashpot (Newton)  $\sigma(t) = \eta \frac{\partial \varepsilon(t)}{\partial t}$

- Linear differential equation or integral equation

$$\sum_{n=0}^N u_n \frac{\partial^n \sigma(t)}{\partial t^n} = \sum_{m=0}^N q_m \frac{\partial^m \varepsilon(t)}{\partial t^m}$$

$u_n, q_m$  : material constants

$$\sigma(t) = \int_0^t G(t-u) \frac{\partial \varepsilon(u)}{\partial u} du$$

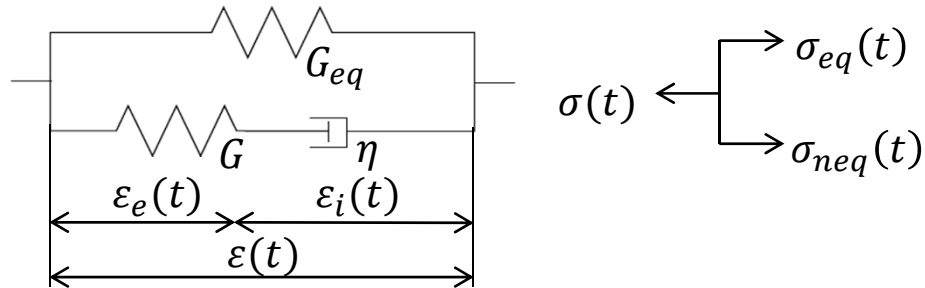
$G(t)$  : relaxation function

- 3D-Formulation

$$\sigma_{ij}(t) = \int_0^t 2G(t-u) \left( \frac{\partial \varepsilon_{ij}(u)}{\partial u} - \frac{1}{3} \frac{\partial \varepsilon_{kk}(u)}{\partial u} \delta_{ij} \right) du + \int_0^t K(t-u) \frac{\partial \varepsilon_{kk}(u)}{\partial u} \delta_{ij} du$$

# Generalized Maxwell Element

- $N = 1$  : Standard linear solid



$$\begin{aligned}\sigma_{eq}(t) &= G_{eq} \varepsilon(t) \\ \sigma_{neq}(t) &= G \varepsilon_e(t) = \eta \frac{\partial \varepsilon_i(t)}{\partial t} \\ \sigma(t) &= \sigma_{eq}(t) + \sigma_{neq}(t) \\ \varepsilon(t) &= \varepsilon_e(t) + \varepsilon_i(t)\end{aligned}$$

$$\frac{\eta}{G} \sigma(t) + \frac{\partial \sigma(t)}{\partial t} = \frac{G_{eq}}{G} \eta \varepsilon(t) + (G + G_{eq}) \frac{\partial \varepsilon(t)}{\partial t}$$

**Linear differential equation**

1. Solution of the homogeneous differential equation  $\sigma_h(t) = e^{\lambda t}$

$$\frac{\eta}{G} \sigma_h(t) + \frac{\partial \sigma_h(t)}{\partial t} = 0 \quad \sigma_h(t) = e^{-\frac{\eta}{G} t} = e^{-\beta t} \quad \beta = \frac{\eta}{G} : \text{decay constant}$$

2. Particular solution: Variation of constants  $\sigma(t) = k(t)e^{-\beta t}$

$$\sigma(t) = \int_0^t (G_{eq} + G e^{-\beta(t-u)}) \frac{\partial \varepsilon(u)}{\partial u} du = \int_0^t G(t-u) \frac{\partial \varepsilon(u)}{\partial u} du$$

**Integral equation**

# Generalized Maxwell Element

- $N \geq 1$  : Prony-Series

$$G(t) = G_{eq} + \sum_{i=1}^N G_i e^{-\beta_i t}$$

$$\sigma(t) = \int_0^t G(t-u) \frac{\partial \varepsilon(u)}{\partial u} du$$

- Linear viscoelasticity for simple, homogeneous and non-aging materials

$$\sigma(t) = \mathcal{F}_{u=0}^{\infty}(\varepsilon(t-u))$$

- Stress-Strain Linearity

$$\begin{aligned} \mathcal{F}_{u=0}^{\infty}(\alpha \varepsilon(t-u)) &= \alpha \mathcal{F}_{u=0}^{\infty}(\varepsilon(t-u)) \\ &= \alpha \sigma(t) \end{aligned}$$

„An increase in the stimulus by an arbitrary factor  $\alpha$  must increase the response by the same factor.“

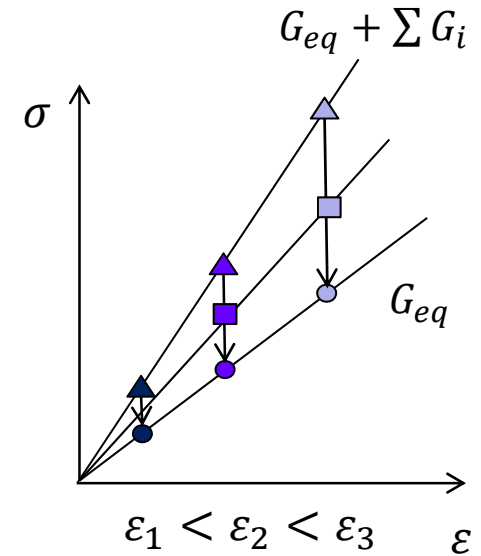
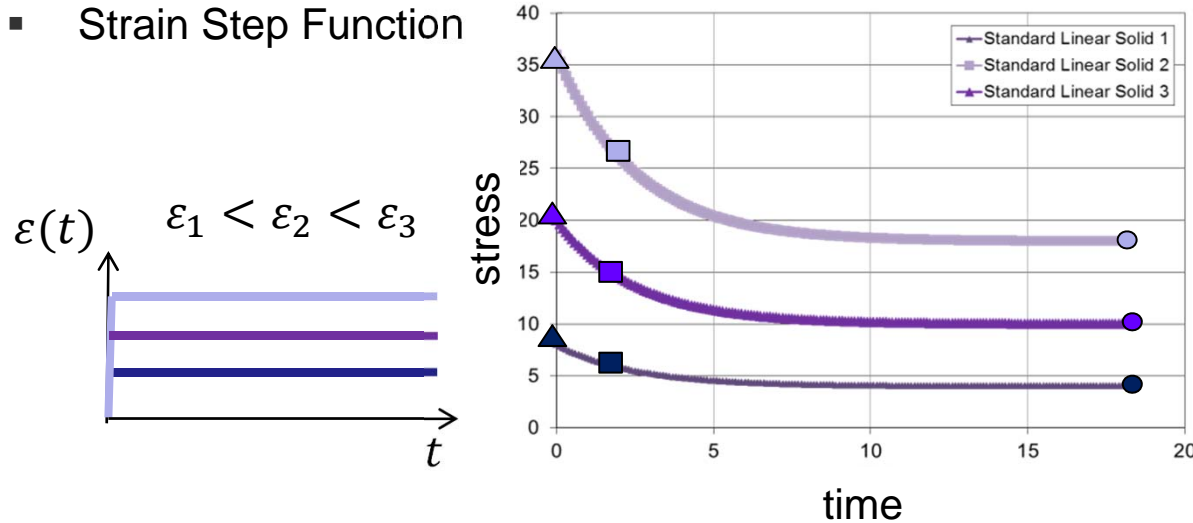
- Boltzmann Superposition Principle

$$\mathcal{F}_{u=0}^{\infty} \left( \sum_{n=1}^{\infty} \varepsilon_n(t-u) \right) = \sum_{n=1}^{\infty} \mathcal{F}_{u=0}^{\infty}(\varepsilon_n(t-u))$$

„An arbitrary sequence of stimuli must elicit a response which is equal to the sum of the responses which would have been obtained if all stimuli had acted independently.“

# Generalized Maxwell Element

- Strain Step Function



$$\varepsilon_1(t) = \varepsilon_1 \mathcal{H}(t)$$

$$\sigma_1(t) = \varepsilon_1 [G_{eq} + \sum G_i e^{-\beta_i t}]$$

Linear equilibrium elasticity  $G_{eq}$

- Fading memory  $\lim_{t \rightarrow \infty} \sigma_1(t) = \varepsilon_1 G_{eq}$

- Stress-Strain Linearity

$$\varepsilon_2(t) = \varepsilon_2 \mathcal{H}(t) = \frac{\varepsilon_2}{\varepsilon_1} \varepsilon_1 \mathcal{H}(t) = \alpha \varepsilon_1(t)$$

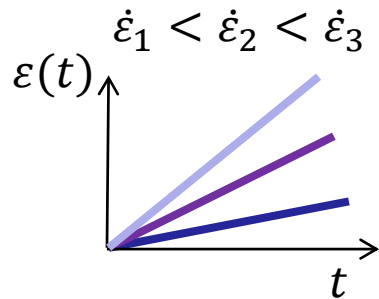
$$\sigma_2(t) = \alpha \sigma_1(t)$$

One test curve defines all other curves.  
Linear viscoelastic material behaviour cannot be identified with one relaxation curve.

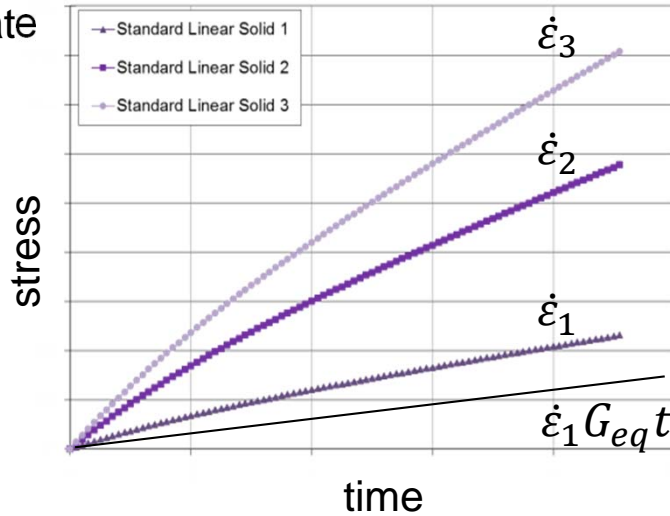


# Generalized Maxwell Element

- Constant strain rate

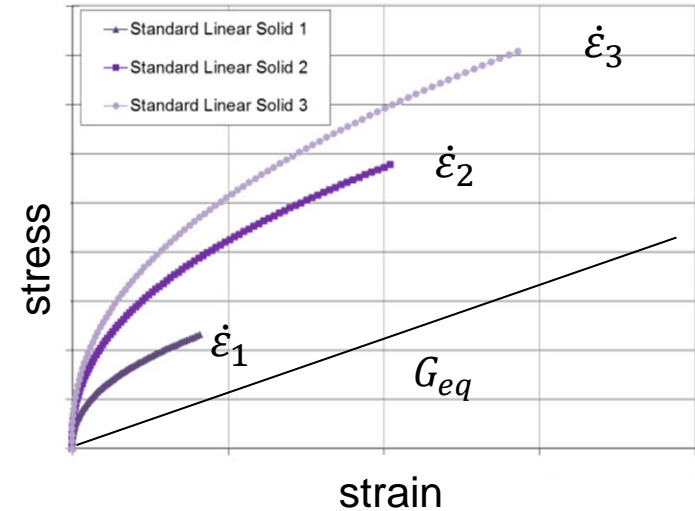


$$\varepsilon_1(t) = \dot{\varepsilon}_1 t$$



$$\sigma_1(t) = \dot{\varepsilon}_1 G_{eq} t + \dot{\varepsilon}_1 \sum \frac{G_i}{\beta_i} (1 - e^{-\beta_i t})$$

$$\lim_{t \rightarrow \infty} \sigma_1(t) = \dot{\varepsilon}_1 G_{eq} t + \dot{\varepsilon}_1 \sum \frac{G_i}{\beta_i}$$



$$\sigma_1(\varepsilon) = G_{eq} \varepsilon + \dot{\varepsilon}_1 \sum \frac{G_i}{\beta_i} (1 - e^{-\beta_i \frac{\varepsilon}{\dot{\varepsilon}_1}})$$

$$\lim_{\varepsilon \rightarrow \infty} \sigma_1(\varepsilon) = G_{eq} \varepsilon + \dot{\varepsilon}_1 \sum \frac{G_i}{\beta_i}$$

- Stress-Strain Linearity

$$\varepsilon_2(t) = \dot{\varepsilon}_2 t = \frac{\dot{\varepsilon}_2}{\dot{\varepsilon}_1} \dot{\varepsilon}_1 t = \alpha \varepsilon_1(t)$$

$$\sigma_2(t) = \alpha \sigma_1(t)$$

Linear viscoelasticity:  
Nonlinear stress-strain curves  
at constant strain rate

# Generalized Maxwell Element

- Equilibrium and non-equilibrium stress

$$\sigma(t) = \mathcal{F}_{u=0}^{\infty}(\varepsilon(t-u))$$

$$\sigma(t) = \int_0^t G(t-u) \frac{\partial \varepsilon(u)}{\partial u} du$$

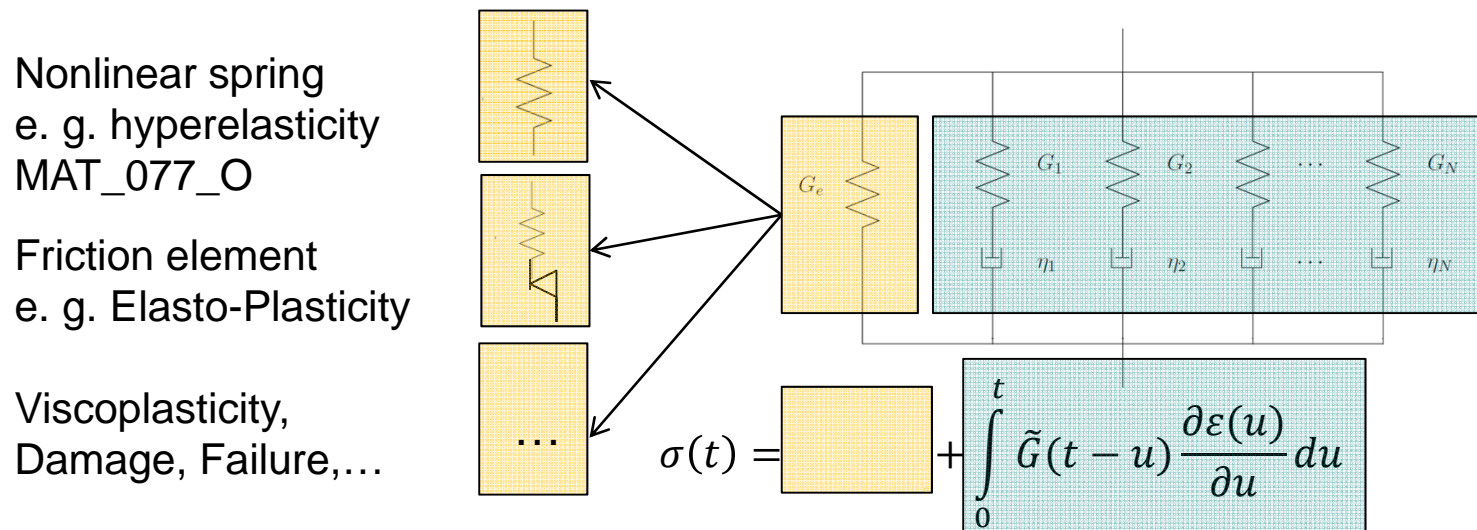
$$G(t) = G_{eq} + \sum_{i=1}^N G_i e^{-\beta_i t}$$

$$\sigma(t) = \mathcal{H}(\varepsilon(t)) + \tilde{\mathcal{H}}_{u=0}^{\infty}(\varepsilon(t-u))$$

$$\sigma(t) = G_{eq} \varepsilon(t) + \int_0^t \tilde{G}(t-u) \frac{\partial \varepsilon(u)}{\partial u} du$$

$$\tilde{G}(t) = \sum_{i=1}^N G_i e^{-\beta_i t}$$

- Non-equilibrium, viscoelastic stress as overstress for different material models

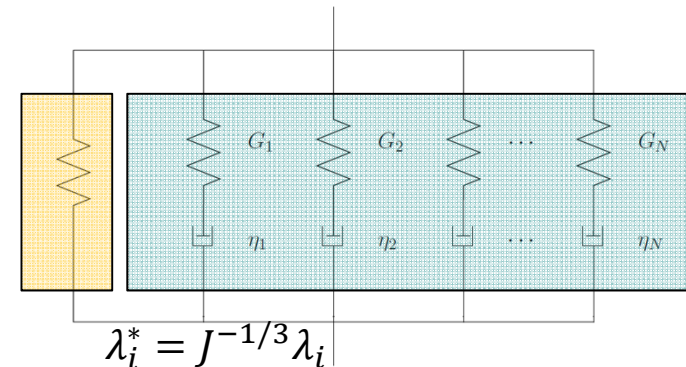


# MAT\_OGDEN\_RUBBER

- Equilibrium stress: Hyperelasticity
  - Ogden strain energy potential and principal stresses

$$W = \sum_{i=1}^3 \sum_{j=1}^m \frac{\mu_j}{\alpha_j} (\lambda_i^{*\alpha_j} - 1) + K(J - 1 - \ln J)$$

$$\sigma_i = \frac{1}{\lambda_k \lambda_j} \frac{\partial W}{\partial \lambda_i} = \sum_{j=1}^m \frac{\mu_j}{J} \left[ \lambda_i^{*\alpha_j} - \sum_{k=1}^3 \frac{\lambda_k^{*\alpha_j}}{3} \right] + K \frac{J-1}{J}$$



$\lambda_i$  : principal stretches  
 $\mu_j, \alpha_j$ : material constants

- Uniaxial loading for incompressible material:  $J \approx 1, \lambda_2 = \lambda_3 = (\lambda_1)^{-1/2} = (\lambda_{uni})^{-1/2}$

$$\sigma_{uni} = \sum_{j=1}^m \left( \mu_j \lambda_{uni}^{\alpha_j} - \mu_j \lambda_{uni}^{-\frac{1}{2}\alpha_j} \right) = \sum_{j=1}^m \left( \mu_j (1 + \varepsilon_{uni})^{\alpha_j} - \mu_j (1 + \varepsilon_{uni})^{-\frac{1}{2}\alpha_j} \right)$$

- Non-equilibrium stress: Viscoelasticity
  - Generalized Maxwell Element for deviatoric deformation

$$\tilde{G}(t) = \sum_{i=1}^N G_i e^{-\beta_i t} \stackrel{v=0.5}{\cong} \sum_{i=1}^N \frac{E_i}{3} e^{-\beta_i t}$$

# Tabulated hyperelasticity

- Hyperelasticity without parameter identification
  - MAT\_FU\_CHANG\_FOAM (MAT\_083)
  - MAT\_SIMPLIFIED\_RUBBER (MAT\_181)

- Ogden Model: Series expansion for incompressible material  $f(\lambda_i^*) = \sum_{j=1}^m \mu_j \lambda_i^{*\alpha_j}$

$$\sigma_i = \sum_{j=1}^m \frac{\mu_j}{J} \left[ \lambda_i^{*\alpha_j} - \sum_{k=1}^3 \frac{\lambda_k^{*\alpha_j}}{3} \right] + K \frac{J-1}{J} = \frac{1}{J} \left[ f(\lambda_i^*) - \frac{1}{3} \sum_{k=1}^3 f(\lambda_k^*) \right] + K \frac{J-1}{J}$$

$$\sigma_{uni} = \sum_{j=1}^m \left( \mu_j \lambda_{uni}^{\alpha_j} - \mu_j \lambda_{uni}^{-\frac{1}{2}\alpha_j} \right)$$

$$f(\lambda_i) = \sum_{j=1}^m \mu_j \lambda_i^{\alpha_j} = \sigma_{uni}(\lambda_i) + \sum_{j=1}^m \mu_j \lambda_i^{-\frac{1}{2}\alpha_j} = \sigma_{uni}(\lambda_i) + \sigma_{uni} \left( \lambda_i^{-\frac{1}{2}} \right) + \sum_{j=1}^m \mu_j \lambda_i^{\frac{1}{4}\alpha_j}$$

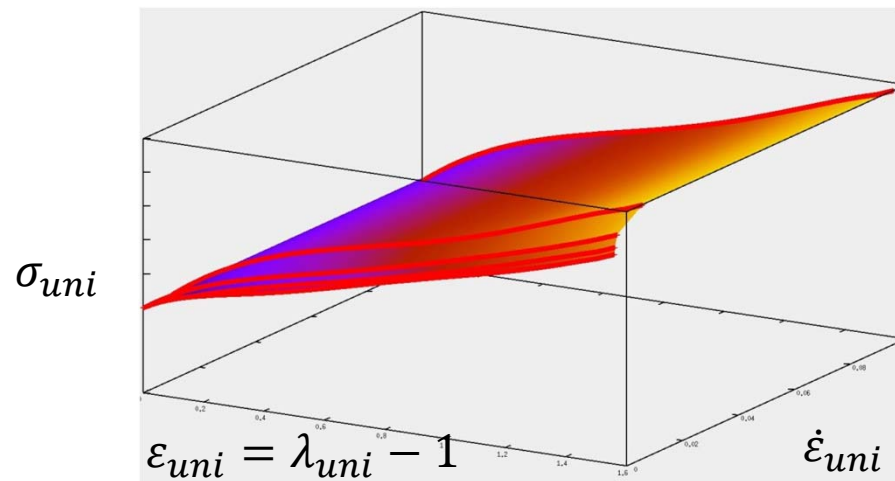
$$= \sigma_{uni}(\lambda_i) + \sigma_{uni} \left( \lambda_i^{-\frac{1}{2}} \right) + \sigma_{uni} \left( \lambda_i^{\frac{1}{4}} \right) + \sum_{j=1}^m \mu_j \lambda_i^{-\frac{1}{8}\alpha_j}$$

$$= \sum_{n=1}^{\infty} \sigma_{uni} \left( \lambda_i^{\left(-\frac{1}{2}\right)^n} \right) \quad \text{Exit if } \left\| \lambda_i^{\left(-\frac{1}{2}\right)^n} \right\| \leq 1.01$$

Principal stress can directly be calculated from uniaxial stress-strain curve.

# Tabulated hyperelasticity

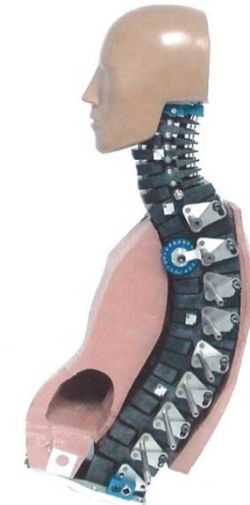
- Rate dependency
  - Loading: Interpolation between strain rates within table definition
  - Unloading: lowest strain rate or damage model



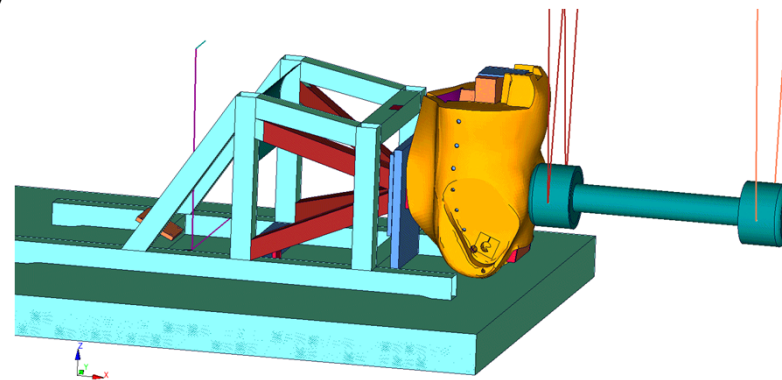
- Damage model
  - Closed loading and unloading path define damage evolution  $d = d(\epsilon_{uni})$
  - Unloading:  $\sigma_{uni} = \sigma_{uni}(1 - d)$

# Modelling of a rubber material

- Silicone – Jacket in BIORID II
- Material characterization
  - incompressible material
  - compression and tensile tests
  - quasistatic and dynamic testing



- Modelling with LS-Dyna: Comparison of strain rate dependent modelling
  - MAT\_SIMPLIFIED\_RUBBER (MAT181)
  - MAT\_OGDEN\_RUBBER (MAT077\_O)
- BioRID Jacket Certification Test
  - Pendulum acceleration
  - Sled acceleration

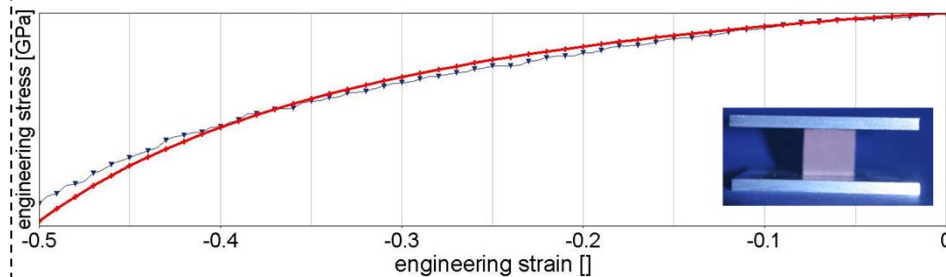
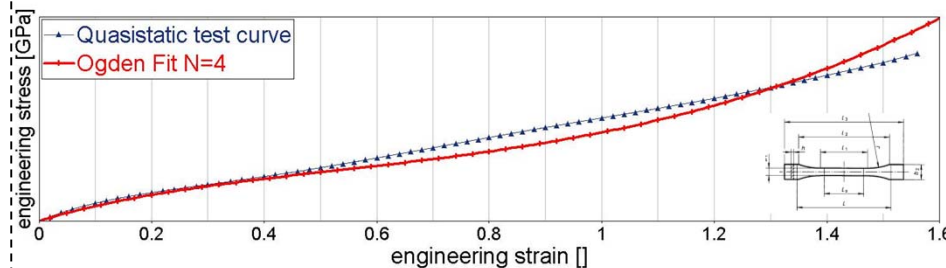


# Material characterization

- Equilibrium stress

Quasistatic test curves –  
Fit to Ogden strain energy potential  
for MAT077 and MAT181

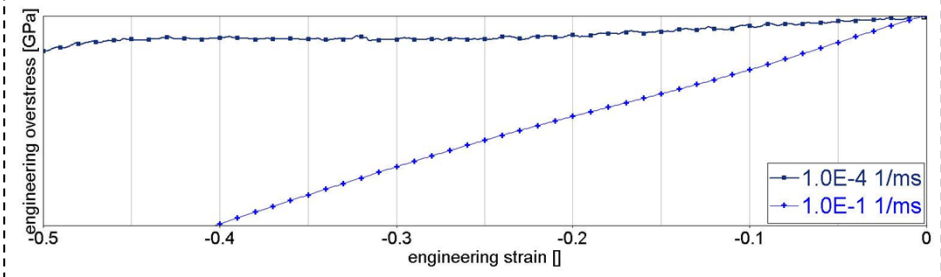
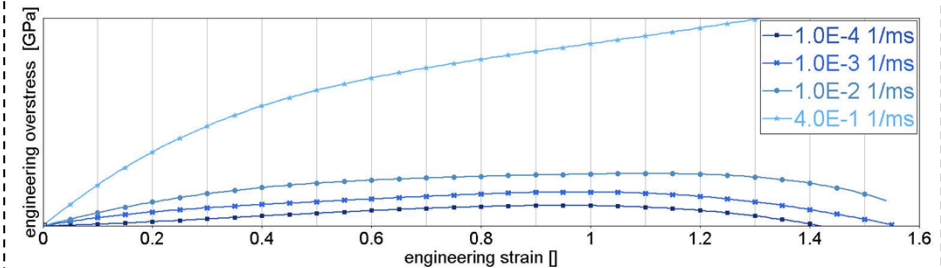
$$\sigma_i = \sum_{j=1}^N \frac{\mu_j}{J} \left[ \lambda_i^{\alpha_j} - \sum_{k=1}^3 \frac{\lambda_k^{\alpha_j}}{3} \right] + K \frac{J-1}{J}$$



- Non-equilibrium stress

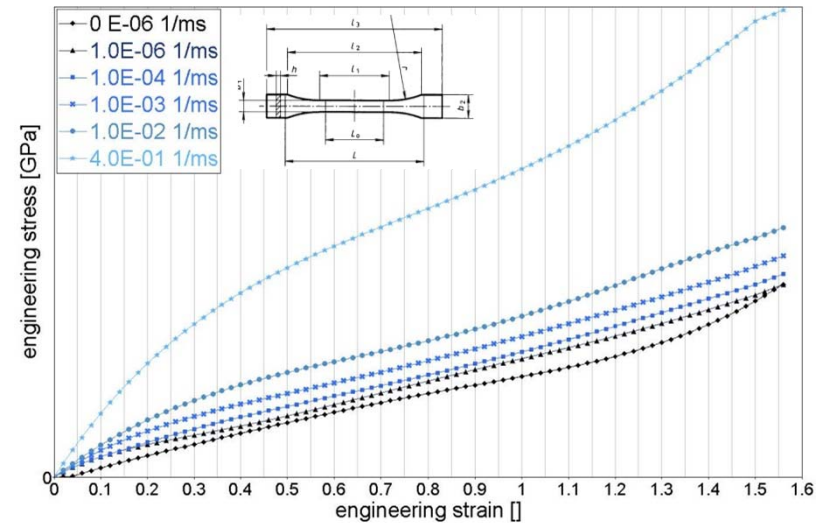
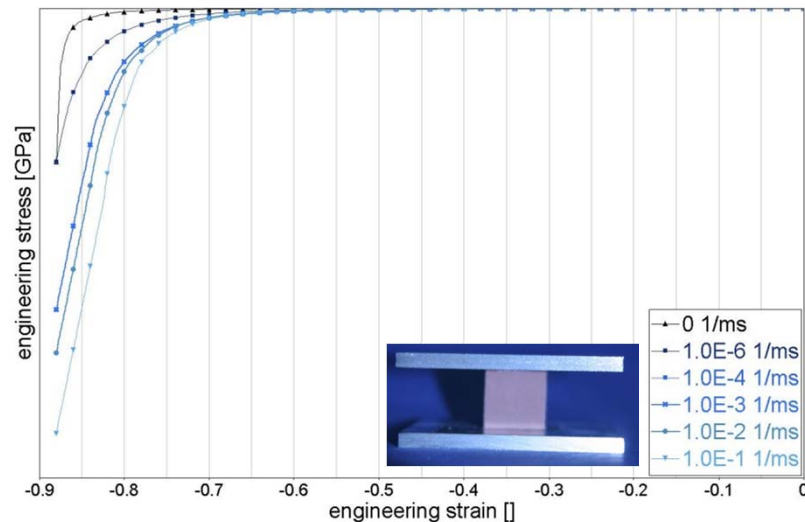
Viscoelastic overstress  
= test curves – Ogden-Fit

- MAT181: Strain-rate dependent table
- MAT077: Generalized Maxwell Element



# MAT\_SIMPLIFIED\_RUBBER

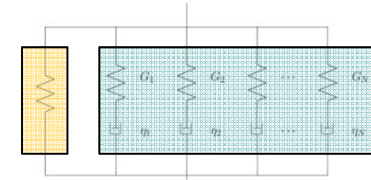
- Table: Ogden-Fit and dynamic test curves



- Strain rate dependency
  - engineering strain rate (RTYPE=1)
  - simple average of 12 time steps (AVGOPT=0)
  - rate effects are treated identically in tension and compression (TENSION=1)
- Unloading
  - Internal damage formulation based on quasistatic unloading path (LCUNLD)
  - Rate dependent unloading path



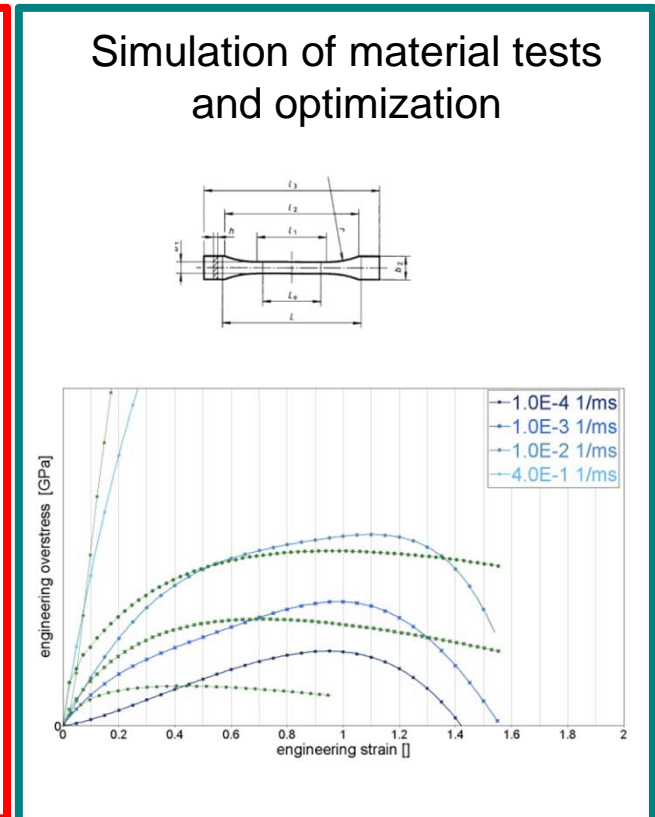
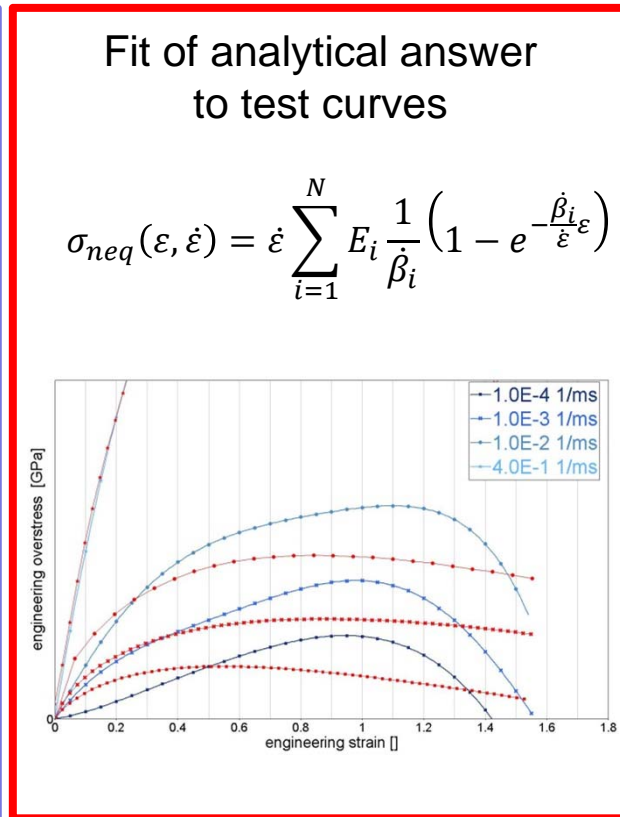
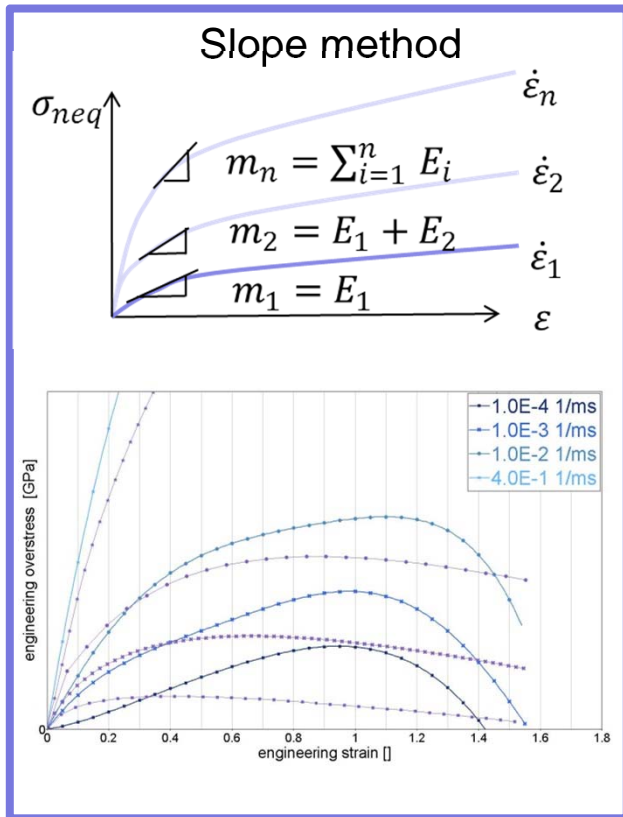
# MAT\_OGDEN\_RUBBER



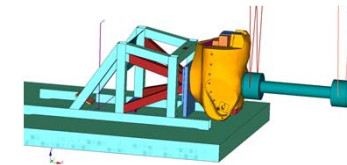
- Parameter identification for Prony-Series
  - uniaxial tensile tests ( $\dot{\epsilon}_i = cst.$ )

$$\tilde{G}(t) = \sum_{i=1}^N G_i e^{-\beta_i t} \stackrel{\nu=0.5}{\cong} \sum_{i=1}^N \frac{E_i}{3} e^{-\beta_i t}$$

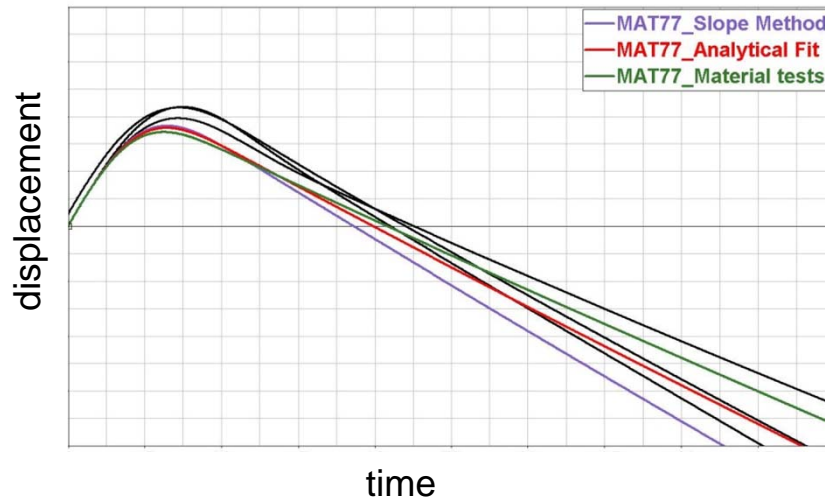
- limited to relevant time scales  $\beta_i = \frac{1}{\dot{\epsilon}_i}$



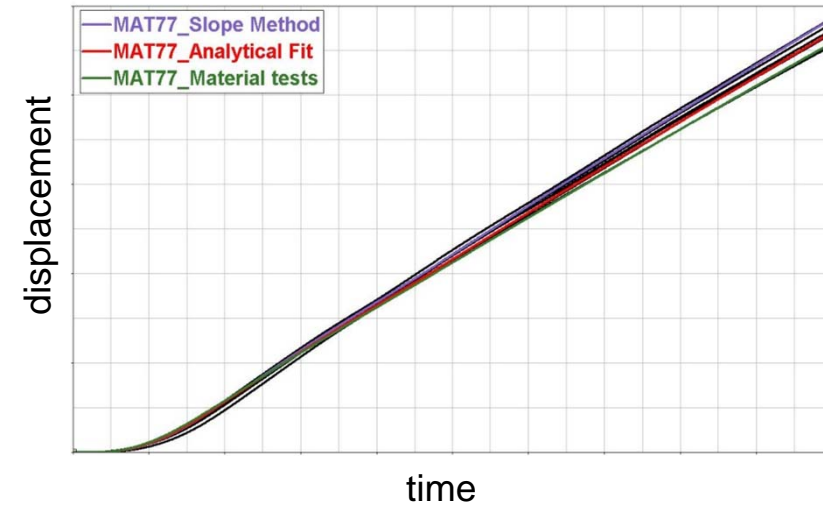
# BioRID Jacket Certification Test: MAT\_77



- Pendulum Displacement

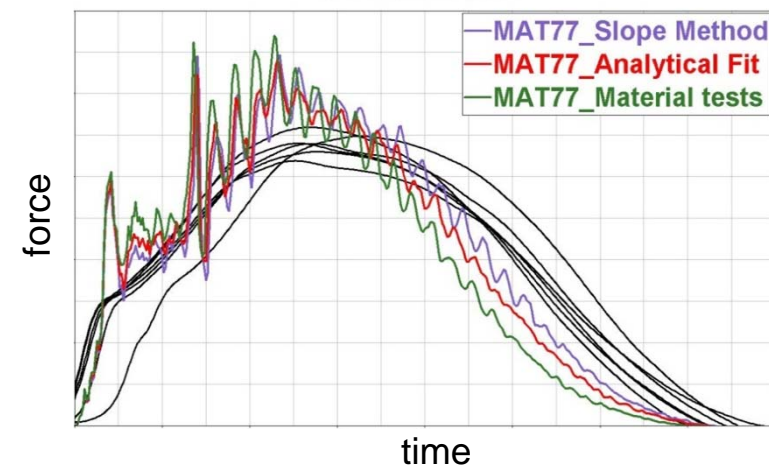


- Sled Displacement



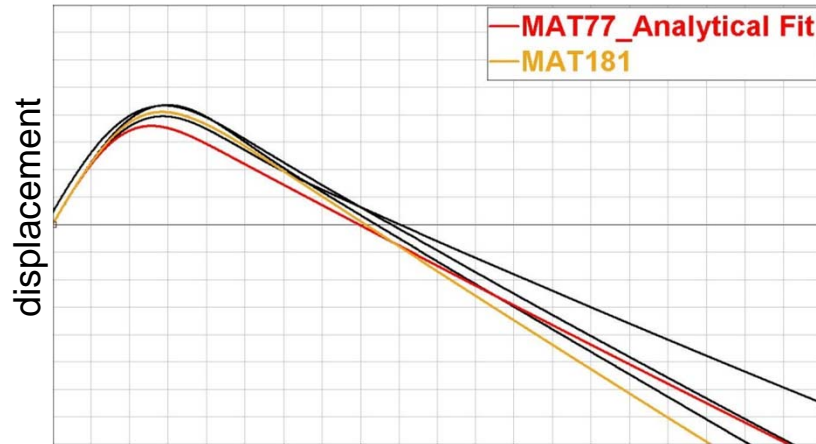
- Minor differences for 3 Prony series
- Initial stiffness: turning point and pendulum force
- Best result for analytical fit

- Pendulum Force

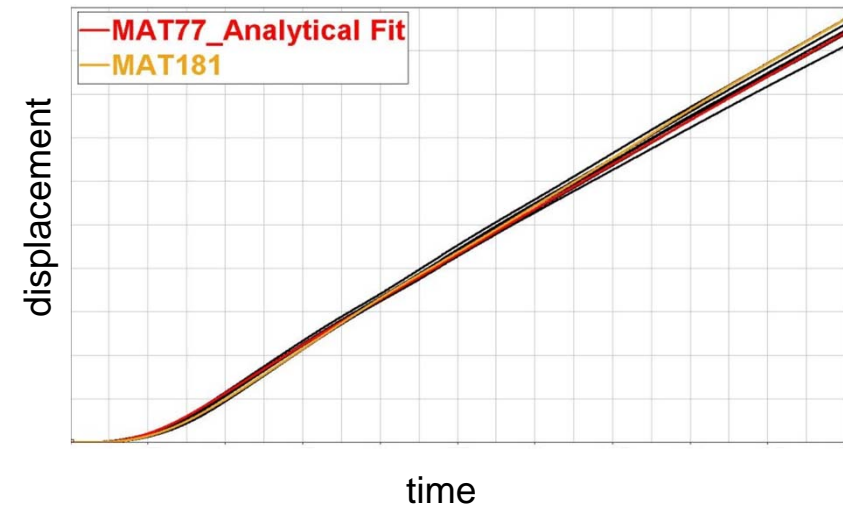


# BioRID Jacket Certification Test: MAT\_77 vs. MAT\_181

- Pendulum Displacement



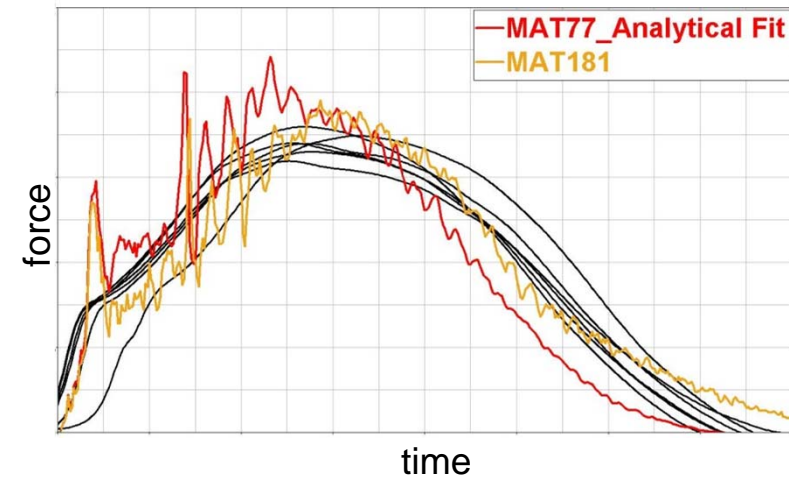
- Sled Displacement



- MAT\_181

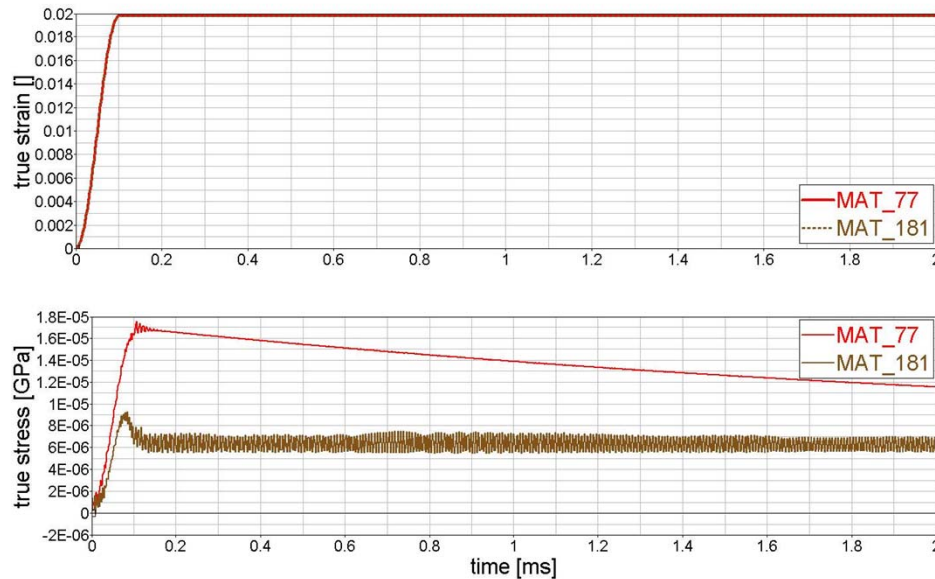
- Better agreement for loading
- Too small hysteresis

- Pendulum Force

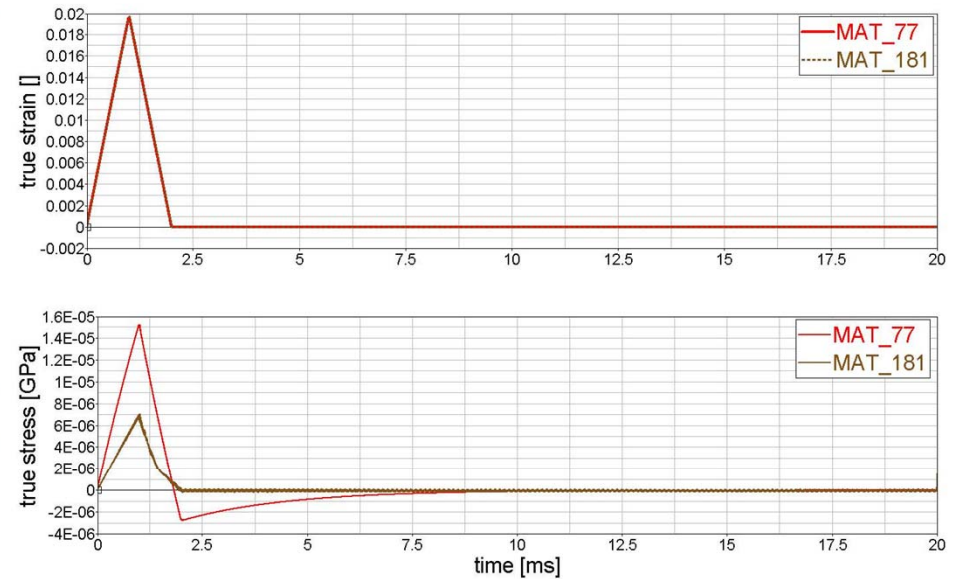


# Tabulated Hyperelasticity vs. Linear Viscoelasticity

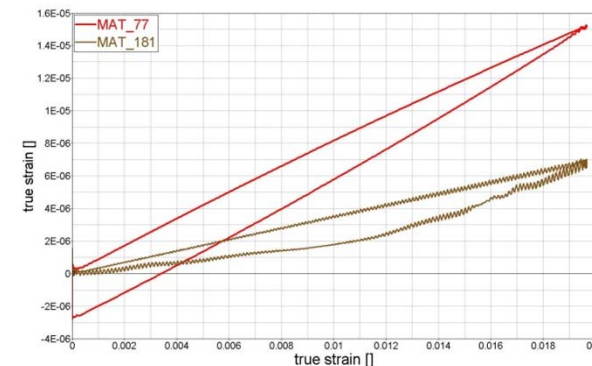
- Relaxation



- Loading and Unloading



- MAT\_181: „Relaxation“ due to averaging over last 12 time steps
- MAT\_77: longer relaxation time leads to sign reversal in the stress upon unloading





# Summary

- Rheological models
  - Linear viscoelasticity
  - No difference between compression and tension
  - Time-consuming parameter identification
  - Test curves differ from model answer
- Tabulated hyperelasticity
  - Nonlinear rate dependency in compression and tension
  - Direct tabulated input using test results
  - Limited modelling of characteristic viscoelastic properties: no creep, no relaxation
  - Difficult to match unloading behaviour in component test simulations



**Thank you for your attention!**